

# §2 – Observational Cosmology – Cosmological theory discussion



<http://www.sr.bham.ac.uk/~smcgee/ObsCosmo/>

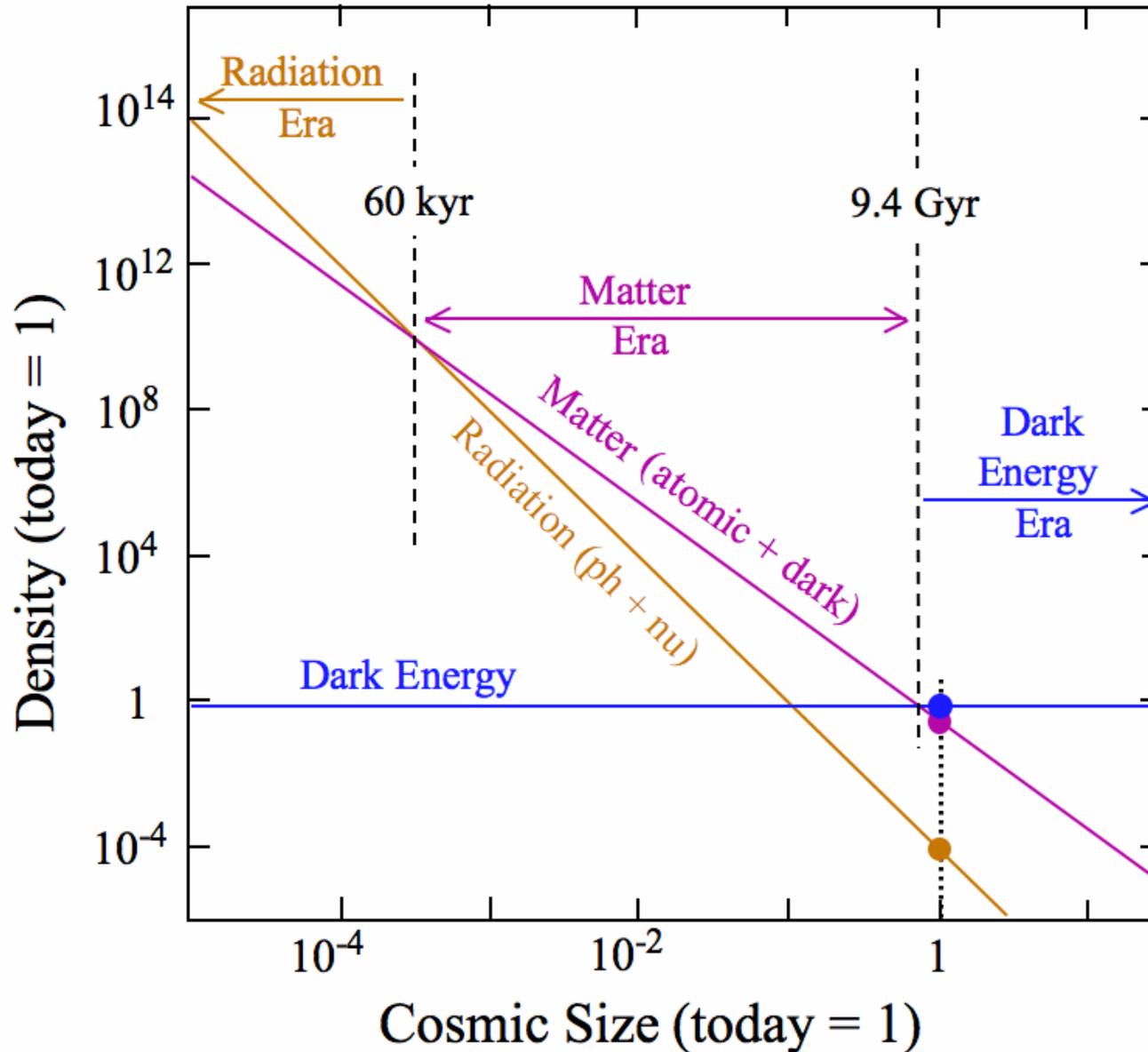
Sean McGee

smcgee@star.sr.bham.ac.uk



Topic	Sources	Comments
<b>Dynamical equations of the Universe:</b> Assumptions - Newtonian treatment The Friedmann equation The acceleration equation	Course lecture , RR(4.3, 5.1), L1(3.0-3.6) L2(3.0-3.6)	
<b>Solutions to the equations with Lambda=0:</b> For matter (dust); radiation; matter + radiation Flat, open and closed Universes Behaviour of the scale factor with time	L1(4), RR(4.6), unit 2 lecture L2(5)	
<b>Cosmological Parameters :</b> H,q,Omega,Lambda - meanings and definitions	L1(6), RR(4.7), unit 2 lecture L2(6)	
<b>The cosmological constant:</b> Friedmann equation with Lambda term Physical interpretation Dynamical solutions - static & accelerating Universes	L2(7), L1(6.4), RR(4.8), unit 2 lecture (& discussion class)	
<b>Curved space-time</b> Flat, spherical and hyperbolic geometries Relation to the Friedmann models	L1(5), L2(4)	
<b>The Robertson-Walker metric</b> The form and meaning of the R-W metric General Relativity also gives the Friedmann equation	RR(4.4, 4.5), L2(A1.1)	
<b>Cosmological Redshift</b> Relation between redshift and expansion (& scale factor)	L1(4.2), RR(7.4) L2(5.2, A2.1)	
<b>Age of the Universe</b> Dependence on Ho, and Omega (for Lambda=0) Comparison with ages of stars Effect of the cosmological constant	L1(7), RR(4.9 & p.163) L2(8)	

# Evolution of the Universe



Evolution of matter, radiation and vacuum contributions

# Critical density

- Friedmann equation

$$\left(\frac{\dot{a}}{a}\right)^2 = H^2 = \frac{8\pi}{3}G\rho - \frac{kc^2}{a^2}$$

- **Critical density** (makes the universe flat,  $k = 0$ )

$$\begin{aligned}\rho_c &= \frac{3H^2}{8\pi G} \\ &= 1.9 \times 10^{-26} h^2 \text{ kg/m}^3\end{aligned}$$

- **Density parameter:**  $\Omega = \frac{\rho}{\rho_c}$

# Density parameter

- Density parameter for “species X”

$$\Omega_X = \frac{\rho_X}{\rho_c}$$

- Curvature density parameter

$$\Omega_k = -\frac{kc^2}{a_0^2 H_0^2}$$

- (A different form for the) Friedmann equation

$$\begin{aligned} H^2 &= \frac{8\pi G}{3} \rho - \frac{kc^2}{a^2} \\ &= \frac{8\pi G}{3} \rho_c \Omega - \frac{kc^2}{a^2} \end{aligned}$$

$$\begin{aligned} \Omega - 1 &= -\frac{kc^2}{a^2} \\ \Omega + \Omega_k &= 1 \end{aligned}$$

# The Einstein Equation(s)

*A tensor equation* - a shorthand for 16 partial differential eqs., connecting the geometry and mass/energy density:

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi GT_{\mu\nu}$$

Spacetime  
geometry

where:

$G_{\mu\nu}$  = The Einstein tensor

$R_{\mu\nu}$  = The Ricci tensor

$g_{\mu\nu}$  = The metric tensor

$R$  = The Ricci scalar

$T_{\mu\nu}$  = The stress-energy tensor

Matter  
distribution

*Homogeneity and isotropy* requirements reduce this set of 16 eqs. to only 1,  $G_{00} = T_{00}$ , which becomes the **Friedmann Equation**

## Metric and Spacetime

Geometry of space can be generally defined through the **metric**, enabling one to compute the distance between any two points:

$$ds^2 = \sum_a \sum_b g_{ab} dx^a dx^b$$

where  $g_{ab}$  is the **metric tensor**. Indices  $\{a,b\}$  run 0 to 3, for the spacetime (0 is the time dimension, 1,2,3 are the spatial ones, i.e., xyz)

In a simple Euclidean geometry, it is a diagonal unit tensor (matrix):  $g_{aa} = 1, g_{a \neq b} = 0$ , where  $\{a,b\} = \{1,2,3\}$

The metric coefficients  $g_{ab}$  are generally functions of the spacetime position, and a proper theory of spacetime has to specify these functions

## Metric: Quantifying the Geometry

- The geometry of spacetime is completely specified by a **metric**,  $g_{\mu\nu}$

- A special relativistic, Euclidean case is the **Minkowski metric**:

$$ds^2 = (c dt)^2 - (dx^2 + dy^2 + dz^2)$$

- A general case for a GR, homogeneous and isotropic universe is the **Robertson-Walker metric**:

$$ds^2 = (c dt)^2 - R^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

where  $k = -1, 0, +1$  for a (negative, flat, positive) curvature

# Robertson-Walker Metric

Polar coordinates are useful if all directions are equal (space is *isotropic*):

$$ds^2 = c^2 dt^2 - R^2(t) \left[ dr^2 + S_k^2(r) d\varphi^2 \right]$$

where

$$S_k(r) = \begin{cases} \sin r, & (k = 1) & \text{Positive space curvature} \\ \sinh r, & (k = -1) & \text{Negative space curvature} \\ r, & (k = 0) & \text{Flat (Euclidean) space} \end{cases}$$

If the spatial dimensions expand or contract with time:

$$ds^2 = c^2 dt^2 - a^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2 d\theta^2 + r^2 \sin\theta d\phi^2 \right]$$

where  $a(t) \equiv \frac{R(t)}{R_0}$  is the **scale factor**

## Introducing the Cosmological Constant

Gravitation is an attractive force, so what is to prevent all matter and energy falling to one gigantic lump?

Einstein introduced a negative potential term to balance the attractive gravity:

$$\nabla^2\phi - \lambda\phi = 4\pi G\rho$$

$\lambda$  could be thought of as an integration constant, or a new constant of nature, or a new aspect of gravity

The Einstein Equations now become:

$$G_{\mu\nu} = 8\pi GT_{\mu\nu} - \Lambda g_{\mu\nu}$$

This is the **cosmological constant**. Note that the theory does not specify its value, or even the sign!

# Essential background cosmology



The evolution of the universe (defined by  $a(t)$ ) is given by the Friedmann equation

$$H^2 \equiv \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G\rho}{3} - \frac{kc^2}{R_0^2 a^2} ,$$

and the acceleration equation

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P/c^2) ,$$

together with the fluid equation and associated equation of state

$$\dot{\rho} = -\frac{3\dot{a}}{a}(\rho + P/c^2) , \quad P = w\rho c^2 .$$

## Matter or Radiation only solutions



Should be able to get the basic matter- and radiation-dominated scalings from the key equations.

Matter:

$$a(t) = \left(\frac{t}{t_0}\right)^{2/3} \quad ; \quad \rho(t) = \frac{\rho_0}{a^3} = \frac{\rho_0 t_0^2}{t^2} \quad . \quad H \equiv \frac{\dot{a}}{a} = \frac{2}{3t} \quad ,$$

Radiation:

$$a(t) = \left(\frac{t}{t_0}\right)^{1/2} \quad ; \quad \rho(t) = \frac{\rho_0}{a^4} = \frac{\rho_0 t_0^2}{t^2} \quad .$$

Radiation-dominated universe expands more slowly because of the deceleration supplied by the pressure term.

## Matter and radiation mixture solutions



Solutions are often usefully obtained if they are matter or radiation-`dominated' but have a mixture of both

Matter-dominated:

$$a(t) \propto t^{2/3} \quad ; \quad \rho_{\text{mat}} \propto \frac{1}{t^2} \quad ; \quad \rho_{\text{rad}} \propto \frac{1}{a^4} \propto \frac{1}{t^{8/3}} .$$

Radiation-dominated:

$$a(t) \propto t^{1/2} \quad ; \quad \rho_{\text{rad}} \propto \frac{1}{t^2} \quad ; \quad \rho_{\text{mat}} \propto \frac{1}{a^3} \propto \frac{1}{t^{3/2}} .$$

Notice the radiation in radiation-dominated case falls off more quickly than matter. So, radiation-dominated case is unstable when there is a mixture.

# Cosmological constant



Friedmann equation with cosmological constant

$$H^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} + \frac{\Lambda}{3}.$$

Acceleration equation with cosmological constant

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\left(\rho + \frac{3p}{c^2}\right) + \frac{\Lambda}{3}.$$

Positive cosmological constant gives a positive contribution to acceleration

# Equation of state of cosmological constant



The cosmological constant can be thought of as a ‘fluid’ with an energy density:

$$\rho_{\Lambda} \equiv \frac{\Lambda}{8\pi G}$$

The fluid equation becomes:

$$\dot{\rho}_{\Lambda} + 3\frac{\dot{a}}{a} \left( \rho_{\Lambda} + \frac{p_{\Lambda}}{c^2} \right) = 0.$$

This must mean:

$$p_{\Lambda} = -\rho_{\Lambda}c^2.$$

More generally:

$$p_Q = w\rho_Qc^2,$$

# Age of the Universe



Estimate from the Hubble constant today:

$$H_0 = 100h \text{ km s}^{-1} \text{ Mpc}^{-1},$$

$$H_0^{-1} = 9.77 h^{-1} \times 10^9 \text{ yrs}.$$

Somewhat more accurately:

$$a(t) = \left( \frac{t}{t_0} \right)^{2/3} \quad \Longrightarrow \quad H \equiv \frac{\dot{a}}{a} = \frac{2}{3t},$$

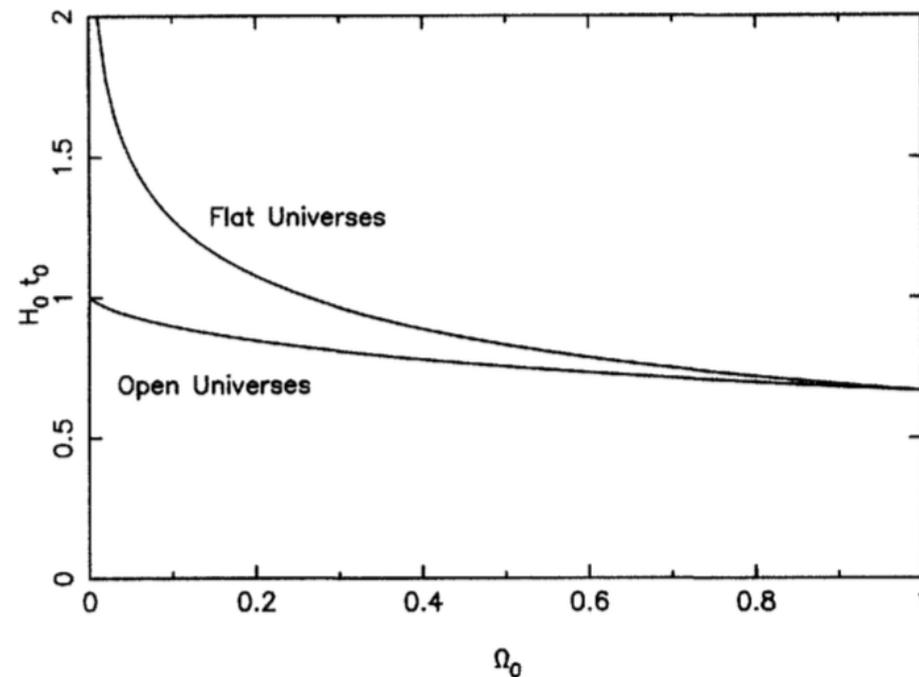
$$t_0 = \frac{2}{3} H_0^{-1} = 6.51 h^{-1} \times 10^9 \text{ yrs}.$$

# Age of the Universe



More complete solution:

$$H_0 t_0 = \frac{2}{3} \frac{1}{\sqrt{1-\Omega_0}} \ln \left[ \frac{1 + \sqrt{1-\Omega_0}}{\sqrt{\Omega_0}} \right] = \frac{2}{3} \frac{1}{\sqrt{1-\Omega_0}} \sinh^{-1} \left[ \sqrt{\frac{1-\Omega_0}{\Omega_0}} \right]$$



# Solutions with a cosmological constant



Now, a 'closed' universe ( $k > 0$ ) doesn't necessarily recollapse.

$$q_0 = \frac{\Omega_0}{2} - \Omega_\Lambda,$$

